

# To cut or not to cut: Effect of vegetation height and bulk density on wildfire propagation under varied wind and slope conditions

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## ABSTRACT

The frequency, intensity and span of wildfires have surged in the past decades, mainly driven by changes in climatic patterns. Although grasslands cover  $\approx 40\%$  of the Earth's surface, they account for  $\approx 80\%$  of the burned area. Mowing is a common practice to reduce the Rate of Spread (RoS) and intensity of grassland fires. However, recent studies suggest the RoS may increase instead. This paper combines results from previous experimental studies in Australian grasslands and wheatlands with wildfire propagation simulations under a range of ambient wind speeds, vegetation heights ( $H_g$ ), and terrain slopes to assess whether *grass-cutting* is an effective strategy to mitigate fire propagation. Previous investigations of the  $H_g$ -RoS correlation led to contradictory conclusions. The main finding here is a statistically significant negative correlation between the vegetation's bulk density ( $\rho_b$ ) and the  $RoS/u_{10}$  ratio, where  $u_{10}$  is the wind speed 10 m above ground. Seeking a relationship between  $H_g$  and RoS, we found positive correlations between  $H_g$  and the  $RoS/(u_{10}M)$  and  $RoS/(u_{10}R_H)$  ratios, where  $M$  (fuel moisture content) and  $R_H$  (atmospheric humidity) account for environmental conditions. These correlations seem to hold provided that  $\rho_b$  decreases with increasing  $H_g$  and that propagation is plume-driven. For wind-driven propagation, simulations show that decreasing  $H_g$  increases  $RoS/u_{10}$  for constant  $M$  and  $R_H$ . Further experimental research is needed to confirm this. Thus, *grass-cutting* is observed to effectively curb fire propagation, although it may be less effective under specific conditions such as high winds, as vegetation characteristics and propagation modes significantly affect fire dynamics.


## 1. Introduction

The complex interactions between biological, climatic, physical, and social factors influence the likelihood of a wildfire ignition, as well as its spread, intensity, duration, and span. Global changes in climate, land use, management techniques, and population are altering wildfire risk in various parts of the world. Areas previously impacted by wildfires might observe changes in risk—whether increasing or decreasing—whereas those that have never had a wildfire before are now at increased risk. Wildfires can potentially devastate roads and other infrastructure, disrupt natural processes such as the supply of water, and cause immediate and long-term adverse effects on public health. In addition, such incidents may interfere with transport and supply chains, leading to road and business closures. Moreover, smoke from wildfires contains hazardous compounds and fine particulates produced by combustion which pose health threats, particularly at the wildland-urban interface (WUI) [45].

Grasslands cover up to  $50M\text{ km}^2$  ( $\approx 37\%$ ) of the Earth's terrestrial surface [36], and comprise more than 80 % of the world's burned land [22]. Grasslands constitute 40 % of land in UK [15] and 70 % in Australia [20]. This highlights the importance of studying grassland fires, which differ from forest fires in various ways. For instance, due to their well-aerated structure, the high surface-to-volume ratio of grass litter, and the lack of trees to obstruct the wind, they may have an extremely high Rate of Spread (RoS), hence posing high levels of danger to people and buildings [42]. The RoS of the fire is generally a function of the interplay between topographical, weather, and fuel factors. These include atmospheric conditions such as ambient wind speed, humidity, and temperature; topographic conditions such as slope; and fuel conditions such as vegetation type, height, density, and moisture content [20, 26].

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*Weather and atmospheric parameters* significantly affect the fire propagation dynamics. Wind, specifically ground-level or near-surface wind, is the most studied parameter in fire–atmosphere interaction. Generally, wind accelerates the RoS of the fire by supplying it with fresh oxygen and by tilting the flame towards the fresh, unburned fuel, leading to an increased preheating of the fuel by diffusion and/or radiation, and transferring the hot air through the wet fuels through a convective process. Additionally, the wind moves firebrands over long distances, causing new outbreaks ahead of the main fire front [40]. Various researchers have investigated the relationship between the *fire RoS* and the *wind speed*. For example, McArthur [23] investigated the effect of the wind on the spread of surface fire in Australian grasslands, proposing that the RoS is a function of the square of the wind speed for winds of up to 10 m/s. However, increasing the wind speed above 12.5 m/s seems to reverse this effect and decrease the RoS instead. One of the most significant contributions to the topic was made by the Commonwealth Scientific and Industrial Research Organisation (CSIRO). In [5], they conducted a large set of experimental grassland fires and examined the effects of wind velocity, fuel moisture content (FMC), fuel load, and fuel height on the behaviour of grassland line fires. They found that even though wind is the most influential factor, other parameters, such as the length of the ignition line, also significantly contribute to the RoS. They also observed two distinct propagation dynamics including (i) a narrow pointed head fire mostly associated with the updraft from the burned area that restricts the lateral spread due to the lateral inflowing winds, and (ii) a broad parabolic-shaped fire with flanks extending beyond the initial ignition length, which tends to propagate faster than the point fire mechanism. In another study [7], combining data from experimental burns and real wildfires, they proposed a new model for predicting the RoS. They proposed a linear relationship between wind speed and RoS for lower wind speeds (below 5 km/h) and a power-law relationship for stronger winds.

Numerous empirical and mathematical models have been developed to predict wildfire behaviour across vegetation types, relying mainly on experimental data and incorporating the Rothermel model [41] and Albini's fuel models [1]. Simple rules of thumb like assuming the RoS to be 10 % [8] or 20 % [9] of the wind speed 10 m above ground level (AGL) are also widely used. These models mainly aid fire services in predicting wildfire RoS during firefighting operations and comprise the basis for more complex models such as FARSITE [14], which provides spatial fire growth estimations. Instead, physics-based models such as the Fire Dynamics Simulator (FDS) [25] and FireProM-F [16, 17] solve the governing equations to simulate fire behaviour with varying levels of detail. A variety of studies have used these models to study wildfire behaviour. For instance, Morvan et al. [29] studied the effect of fire intensity and wind conditions on the unsteady behaviour of the fire front via two-dimensional (2D) simulations. Analysing temporal variations in fire intensity, they found that plume-dominated flames oscillate more rapidly, exhibiting erratic behaviour that is less predictable than that of wind-driven fires.

*The terrain slope's effect* on the fire behaviour has also been extensively studied, often in combination with wind speed. Rothermel [41] developed a mathematical model that describes the relationship between fire RoS with mid-flame height wind speed and terrain slope. The model quantifies additional propagating flux produced by wind and slope by defining wind and slope coefficients:  $\phi_w$  and  $\phi_s$ , respectively. These coefficients contribute to the overall RoS by being combined with the RoS in flat and windless conditions ( $RoS_0$ ) as follows:  $RoS = RoS_0 (1 + \phi_w + \phi_s)$ . Weise et al. [47] found that increasing wind speed increases both RoS and flame length. They also validated existing mathematical models against experimental data, and concluded that the existing formulations of empirical fire spread models are inaccurate and need to be revised. Wu et al. [48] studied the interaction of a pool-fire plume with the terrain slope under no-wind conditions, observing that flame bends toward the surface and attaches to the bed as the terrain slope increases. They concluded that this is mainly due to the asymmetric formation of the plumes on either side of the flame. Morandini et al. [28] studied the fire spread dynamics uphill under no-wind conditions using particle image velocimetry and video imaging. They observed that, on horizontal surfaces with radiation as the dominant preheating mechanism, the fire plume maintains a quasi-vertical shape due to the lateral air flow into the fire from either side. By increasing the terrain slope, a strong convective flow forms ahead of the flame due to the pressure difference upstream and downstream of the flame that blows towards the top of the surface contributing to the fuel preheating. Monroy et al. [44] studied the effect of the fuel depth and packing ratio on the fire propagation dynamics at different uphill angles using numerical simulations. Consistently with other studies, they identified a critical slope angle of approximately  $22^\circ$  beyond which rapid increase of the RoS occurs with increasing slope.

Although most studies on upslope fire propagation focus on no-wind conditions, a few have examined the combined effects of wind and slope. Pimont et al. [38] examined the integrated effect of wind and slope on fire RoS, emphasising the impact of fire width using the HIGRAD-FIRETEC model. Their results indicate a significant interaction between wind and slope so that, under low-wind conditions, the RoS increases exponentially with slope until the latter reaches

40 %. Instead, the RoS and slope are linearly related under high-wind conditions. Guo et al. [19] experimentally examined the effects of slope ( $0^{\circ}$ – $30^{\circ}$ ) and wind (0–2 m/s) on surface fire spread over a pine needle fuel bed. Their results indicate that the RoS, flame length, and heat flux increase with increasing wind speed and slope, whilst the flame angle decreases. Additionally, they concluded that the RoS increases linearly with slope for low wind speeds  $u < 0.5$  m/s and slopes of up to  $30^{\circ}$ . However, the RoS accelerates abruptly for higher wind speeds as the slope exceeds  $25^{\circ}$ , constituting an example of extreme and eruptive wildfire behaviour.

*The fuel characteristics's effect* on the fire dynamics is controversial, as the literature reports contrasting results. Results from CSIRO experiments [5] indicate that, even though fuel load and vegetation species do not significantly affect fire spread, fires in natural undisturbed vegetation burn approximately 18 % faster than those in cut or grazed ones. Moinuddin et al. [26] utilised FDS to study the effect of relative humidity and fuel moisture content (FMC) on grass fire dynamics. They concluded that reduced humidity and FMC increase RoS and burn intensity, with potential shifts in fire propagation modes. Our study only focuses on vegetation height and bulk density. Cruz et al. [13] conducted 58 experimental fires to study the effect of fuel load on fire behaviour in Australian grasslands. They reported that contrary to the community's common assumptions, a negative correlation between the fuel load and fuel height with the RoS of fire is observed when fuel load is not a limiting factor. In a related study, Moinuddin et al. [27] investigated the effect of fuel height on the fire RoS considering different heights of up to 0.6 m. They concluded that increasing the grass height while keeping a constant bulk density reduces the fire RoS but increases its Intensity ( $I$ ) and Heat Release Rate (HRR), also shifting the fire propagation mode from wind-driven to plume-dominated. However, Cruz et al. [11] later studied the effect of fuel characteristics on the fire RoS in wheat farms by carrying out 45 experimental burns, finding a positive correlation between vegetation height and RoS, where unharvested grass yielded the fastest RoS and longest flames. Commenting on the findings in [13, 27], Cruz et al. [12] stated that the proposed conclusions are counter-intuitive, as robust empirical evidence supports the positive correlation between fuel height and fire RoS. They reanalysed the data from [5, 11] and stated that the positive correlation of fire RoS with vegetation height is not a matter of debate, arguing that variations in vegetation structure may explain why some studies like [13] fail to observe this relationship. In response to these comments, Sutherland et al. [43] reanalysed the data in [11] combined with a series of numerical simulations (with constant  $\rho_b$ ) and stated that, below a certain grass height ( $\approx 0.2$ – $0.24$  m), RoS increases with height, which is consistent with [5, 11], where most fires were wind-driven. However, for taller grasses where fires are typically plume-dominated, the RoS decreases with increasing grass height (negative correlation).

Thus, the effect of grass height on fire RoS is still an open question. Previous studies have reported conflicting results, whilst numerical studies rely on simulations with a constant bulk density—which is not the case in natural vegetation. This becomes even more complicated on sloped terrains, where the combined effect of terrain slope, wind conditions, and vegetation height can lead to unexpected results, which have not yet been adequately studied in the literature. Consequently, this paper attempts to address this gap by performing a series of numerical simulations and combining the results with two contradictory experimental studies to investigate the effect of vegetation height and structure on fire propagation dynamics for different terrain slopes and ambient wind speeds. A total of 84 field scale simulations are conducted, and the results are analysed and discussed with regards to the effect of each parameter on the fire RoS and propagation modes.

The remainder of this paper is organised as follows: Section 2 provides an overview of the fire propagation model used in this research (Section 2.1), the reference case and fuel definitions (Section 2.4), the turbulent wind model (Section 2.5), the model gridding (Section 2.6), and the model reliability (Section 2.7); Section 3 presents the results of simulations, including the fire RoS under different wind (Section 3.1) and terrain (Section 3.2) conditions, and the characterisation of the effect of vegetation height and bulk density on the fire dynamics; Section 4 provides a discussion of the results, including an attempt to answer the question of whether *grass cutting* is an adequate management strategy to curb the propagation of the fire; whilst Section 5 provides a summary of the research findings and derived conclusions.

## 2. Materials and Methods

### 2.1. Fire dynamics simulator

The Fire Dynamics Simulator (FDS) numerically solves a version of the Navier-Stokes equations adapted to low-speed and thermally-driven flows, emphasising the simulation of smoke and heat transmission from flames. The core algorithm uses an explicit predictor-corrector technique with second-order accuracy in both space and time.

Large Eddy Simulation (LES) is used to model turbulence within the solution domain. However, for sufficiently fine meshes, Direct Numerical Simulation (DNS) can be employed as an alternative. FDS typically employs a one-step, mixing-controlled chemical reaction model involving three bundled species: *products*, *fuel*, and *air*. Under certain circumstances, reactions that are not always mixing-controlled and multiple reactions might be taken into consideration. A Cartesian grid is used to discretise the domain and estimate the governing equations [24]. FDS offers different models for stimulating wildfire spread depending on the level of physical details required and the computational resources available: the Lagrangian Particle Model (LPM), the Boundary Fuel Model (BFM), and the Level-Set (LS) model, which uses the same elliptical spread model as FARSITE based on Huygens' principle for wave-front propagation modelling, and adopts Rothermel-Albini's RoS formula and Albini's 13 fuel models.

### 2.1.1. Boundary fuel model

This model may be used when a coarse grid is desired to discretise a thin layer of vegetation. Here, the vegetation is represented as a porous barrier made up of a layer of wetness, air, and dry vegetation. For grid sizes up to 10 m, the vegetation height can be used although it is not resolved on the grid [24, 46]. In this model, the convective heat transfer is represented by a source term in the one-dimensional heat conduction equation. This equation is applied to both the vegetation layer and the solid ground. Additionally, the transfer of thermal radiation through the vegetation layer is modelled using a one-dimensional radiative transport equation designed for semi-transparent solids. Even though this model might be efficiently utilised in a variety of studies and simulations, our preliminary results indicate that the LPM provides better and more realistic results. Consequently, this study utilises the LPM for the simulations.

### 2.1.2. Lagrangian particle model

In this model, a group of Lagrangian particles heated by convection-radiation heat transfer represents the vegetation. These particles may be grass, trees, leaves, or anything else. LPM may be used to replicate the front, rear, and flank fire across the surface and high-level vegetation (e.g. trees) with appropriate grid refinement [46]. The drag force per unit volume ( $\mathbf{f}_b$ ) exerted by the vegetation is modelled as follows:

$$\mathbf{f}_b = \frac{\rho}{2} C_d C_s \beta \sigma \mathbf{u} \|\mathbf{u}\| \quad (1)$$

where  $\rho$  is the air density,  $C_d$  is the drag coefficient defined through laboratory experiments,  $C_s$  is the shape factor with the default value of 0.25,  $\beta$  is the packing ratio of vegetation calculated as mass per unit volume divided by material density,  $\sigma$  is the surface-area-to-volume ratio, and  $\mathbf{u}$  is the wind velocity.

## 2.2. Dimensional analysis

Performing dimensional analysis using the Buckingham-pi theorem, Morvan et al. [30] stated that the propagation of fire in grasslands could be described by three non-dimensional parameters, namely ( $RoS/u_w$ ), Byram's convective number ( $N_c$ ), and the fuel moisture content ( $M$ ), where

$$\frac{RoS}{u_w} = F(N_c, M), \quad (2)$$

$$N_c = \frac{2gI}{\rho c_p T_0 (u_w - RoS)^3}, \quad (3)$$

$I$  represents fire intensity,  $\rho = 1.225 \text{ kg/m}^3$  is the air density,  $c_p = 1010 \text{ J/kg/K}$  is the specific heat capacity, and  $T_0$  is the ambient temperature in Kelvin. Fire intensity  $I = W \times H \times RoS$  (kW/m),  $W$  (kg/m<sup>2</sup>) is the fuel load, and  $H$  (kJ/kg) is the heat of combustion of the fuel. Nelson [33] suggested that for values of  $N_c < 2$ , the propagation of fire is mainly dominated by the convective heat transfer between the flame and the unburned vegetation ahead of it: the wind-driven fire propagation mode. Instead,  $N_c > 10$  leads to a different fire propagation mode, characterised by a vertical visible plume that is mostly governed by buoyancy forces and radiative heat transfer between fuel and flame. Consequently,  $2 < N_c < 10$  may be considered a transitory state of fire. While  $u_w$  is often used interchangeably with both  $u_2$  and  $u_{10}$  in the literature when calculating  $N_c$ , this study uses  $u_{10}$  in all calculations for consistency.

**Table 1**

Measured properties of CSIRO C064 and F19 experiments [5, 25].

Property	Unit	Case C064	Case F19 quad
Wind speed ( $u_2$ )	m/s	4.6	4.8
Ambient Temperature ( $T$ )	°C	32	34
Surface Area to Volume Ratio ( $\sigma$ )	m <sup>-1</sup>	9,770	12,240
Grass Height ( $H_g$ )	m	0.21	0.51
Bulk Mass Per Unit Area ( $\rho_b$ )	kg m <sup>-2</sup>	0.283	0.313
Moisture Fraction ( $M$ )	%	6.3	5.8
Measured $RoS$	m s <sup>-1</sup>	1.2	1.5

### 2.3. Reference Cases

Numerical simulations are validated against the field scale grass fire propagation experiments carried out in the Commonwealth Scientific and Industrial Research Organisation (CSIRO) fields during July and August 1986 with constant high daily temperatures [5, 7]. During the experiments, air temperature, relative humidity, and solar radiation were measured 1.4 m above ground level (AGL), while wind velocity was measured 2 m AGL. Fuel load and height were sampled at 16 points for each experiment and averaged for each plot. Other fuel characteristics, including surface-area-to-volume ratio and the fuel moisture content (FMC), were measured and recorded for each experiment. Measured properties of two experiments from [5] are provided in Table 1. The C064 experiment was carried out in a 100×100 m<sup>2</sup> field covered with dry Kerosene grass of 0.21 m high, while the F19 experiment was carried out in a 200×200 m<sup>2</sup> field covered with dry Kangaroo grass of 0.51 m high. Two workers ignited fires, starting at the midpoint and moving towards the sides of the plot. The length of ignition was 50 m for the C064 and 175 m for the F19 experiments. The fire behaviour and RoS were studied using data gathered from ground observations and oblique aerial photographs.

### 2.4. Fuel Modelling

Following [31], the solid phase thermal degradation of the vegetation is modelled using a three-step reaction process including (i) endothermic moisture evaporation, (ii) endothermic pyrolysis of dry vegetation, and (iii) exothermic char oxidation. Thus, the rate of change of the total mass in terms of density of composite solid is calculated as:

$$\frac{\partial \rho_s}{\partial t} = -r_{H_2O} - (1 - v_{char}) r_{pyr} - (1 - v_{ash}) r_{char} \quad (4)$$

where the reaction rates for evaporation of H<sub>2</sub>O, pyrolysis of the dry vegetation, and surface oxidation of char as a function of component densities of composite solid are calculated via Arrhenius kinetics:

$$r_{H_2O} = \rho_{s,H_2O} A_{H_2O} T^{-\frac{1}{2}} e^{\left(-\frac{E_{H_2O}}{R T}\right)} \quad (5)$$

$$r_{pyr} = \rho_{s,dry} A_{pyr} e^{\left(-\frac{E_{pyr}}{R T}\right)} \quad (6)$$

$$r_{char} = Y_{O_2,surf} \sigma A_{char} e^{\left(-\frac{E_{char}}{R T}\right)} \quad (7)$$

where  $\sigma$  is the surface area-to-volume ratio of the vegetation,  $Y_{O_2,surf}$  is the oxygen mass fraction at the material surface,  $A$  and  $E$  (J/mol) are the pre-exponential factors and activation energy,  $R = 8.314$  J/mol/K is the universal gas constant, and  $T$  is the absolute temperature in Kelvin [24]. This approach allows for the possibility of parallel drying and pyrolysis, with series char oxidation taking place as char is being produced during pyrolysis [31]. Kinetic constants must be determined for this model, and as the data required for particular fuels is sometimes unavailable—including the vegetations in CSIRO experiments—the kinetic constants of Fir determined by Grishin [18] are used for the pyrolysis model in this study. Table 2 presents the details of the pyrolysis constants, as well as the soil and two vegetation models utilised in this study. It is important to note that  $Veg_1$  is defined so as to reproduce the Kerosene grass utilised in the CSIRO C064 experiment, while  $Veg_2$  is defined with the lower bulk density and the higher area-to-volume ratio of Kangaroo grass in the F19 experiment.

**Table 2**

Summary of fuel physical properties and thermal decomposition coefficients.

Property	Unit	Veg <sub>1</sub> (Veg <sub>2</sub> )	Reference
Area to Volume Ratio ( $\sigma$ )	m <sup>-1</sup>	9,770 (12,240)	[5]
Bulk density ( $\rho_b$ )	kg.m <sup>-3</sup>	1.313 (0.616)	[5]
Fuel	-	Cellulose	[46]
Fuel Density ( $\rho$ )	kg.m <sup>-3</sup>	512	[46]
Moisture content ( $M$ )	%	6.3	[5]
Specific Heat	kJ kg <sup>-1</sup> K <sup>-1</sup>	2.1	[2]
Conductivity	kJ kg <sup>-1</sup> K <sup>-1</sup>	0.1	[46]
Heat of Evaporation ( $H_{H_2O}$ )	kJ kg <sup>-1</sup>	2259	[39]
Heat of Combustion ( $H_c$ )	kJ kg <sup>-1</sup>	17,400	[24]
Heat of Pyrolysis ( $H_{pyr}$ )	kJ kg <sup>-1</sup>	418	[39]
$A_{pyr}$	s <sup>-1</sup>	1040	[18]
$E_{pyr}$	J.mol <sup>-1</sup>	61041	[18]
Char Yield ( $v_{char}$ )	kg kg <sup>-1</sup>	0.25	[24]
$A_{char}$	kg.m <sup>-2</sup> .s <sup>-1</sup>	465	[3]
$E_{char}$	J.mol <sup>-1</sup>	68000	[3]
Ash Yield ( $v_{ash}$ )	kg kg <sup>-1</sup>	0.04	[24]
Obukhov Length ( $L$ )	m	-500	[46]
Roughness Length( $z_0$ )	m	0.03	[46]
Drag Coefficient ( $c_d$ )	-	2.8	[24]
Soil Specific Heat	kJ kg <sup>-1</sup> K <sup>-1</sup>	2.0	[46]
Soil Conductivity	W m <sup>-1</sup> K <sup>-1</sup>	0.25	[46]
Soil Density	kg m <sup>-3</sup>	1,300	[46]
Relative Humidity	%	40	[46]

## 2.5. Turbulent wind model and boundary conditions

FDS offers four options for defining the inlet wind into the domain, including (i) a specified wind speed and direction that remains constant with height, (ii) the Monin-Obukhov similarity theory, (iii) advanced meteorological concepts such as a Geostrophic wind for modelling huge spatial domains, and (iv) the power law approximation or the "wall of wind" model. The latter is widely used due to its simplicity, mainly with a power of 1/7 [25, 20, 27]. However, this model has limitations, as the power value varies with height, surface roughness, and stability conditions [37]. The 1/7 exponent approximates a neutrally stable atmospheric stratification. Notably, the CSIRO C064 experiments were carried out during summer daytimes under consistently warm and dry conditions ( $T_{C064} = 32^\circ\text{C}$ ) [5], resulting in an unstable stratification of the atmosphere characterised by convective uprising of warm surface air [35]. Thus, the *Monin-Obukhov* similarity theory, which estimates the vertical wind and temperature profiles based on surface and atmospheric conditions [24], would probably provide a more accurate representation. This theory assumes that wind speed ( $u$ ) and potential temperature ( $\theta$ ) change with height as follows [24]:

$$u(z) = \frac{u_*}{k} \left[ \ln \left( \frac{z}{z_0} \right) - \Psi_m \left( \frac{z}{L} \right) \right] \quad (8)$$

$$\theta(z) = \theta_0 + \frac{\theta_*}{k} \left[ \ln \left( \frac{z}{z_0} \right) - \Psi_h \left( \frac{z}{L} \right) \right] \quad (9)$$

where  $u_*$  is the friction velocity,  $k = 0.41$  is the von Karman constant,  $z_0$  is the aerodynamic roughness length,  $\theta_*$  is the scaling potential temperature,  $\theta_0$  is the ground level temperature,  $L$  is the Obukhov length, and  $\Psi_h$  and  $\Psi_m$  represent similarity functions. A negative value of  $L$  (m) determines an unstable stratified atmosphere where the buoyancy-generated turbulence causes large fluctuations in wind velocity and direction and enhances mixing. Based on the suggestions in [46], this study determines an Obukhov length of  $L = -500$  m and a roughness length of  $z_0 = 0.03$  m.



Turbulence within the domain is simulated utilising the very large eddy simulation (VLES) model, incorporating Deardorff's sub-grid scale (SGS) model to handle turbulent eddy viscosity closure terms. The Van Driest damping model is applied to model the Reynolds stresses in near-wall regions. To replicate the turbulent nature of natural atmospheric winds which significantly impact their behaviour at domain boundaries, this study employs the synthetic eddy method (SEM) [21]. This introduces random eddies into the domain, as in [20, 32]. Given that an accurate representation of the eddy characteristics requires measurements of the turbulent Reynolds stress within the ambient wind and canopy height (see [32]), which are not available here, an arbitrary turbulence intensity of 10% is determined (see [20]). It should be noted that the wind field is allowed to develop throughout the simulation domain for 80 s before igniting the fire. Boundary Conditions include a no-slip condition for the ground surface, with 'open' boundary conditions applied to the rest of the boundaries.

## 2.6. Simulation domain and gridding

The burnable field simulated in this study includes a vegetation field of  $200 \times 200 \text{ m}^2$ , allowing sufficient time and space for the fire to reach quasi-steady state, particularly under steep upslope and tall vegetation conditions. Preliminary experiments indicated that increasing domain height beyond 40 m does not affect simulation results. Additionally, in order to eliminate the effects of upwind and downwind boundaries on the fire propagation, a minimum buffer of 150 m was set on both sides. Consequently, the simulation domain of this study, as shown in Fig. 1, is a cuboid measuring 600 m in length, 320 m in width, and 60 m in height.

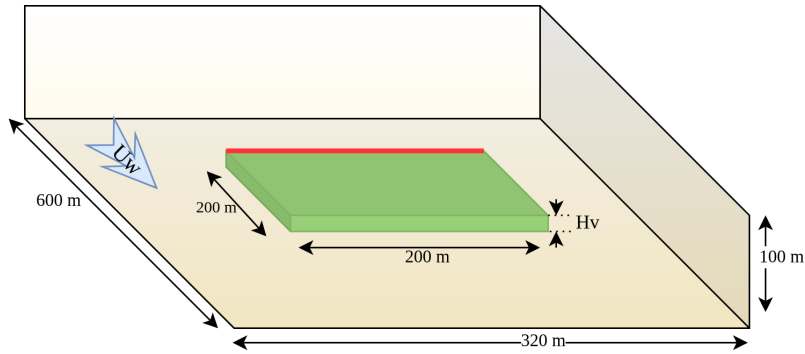


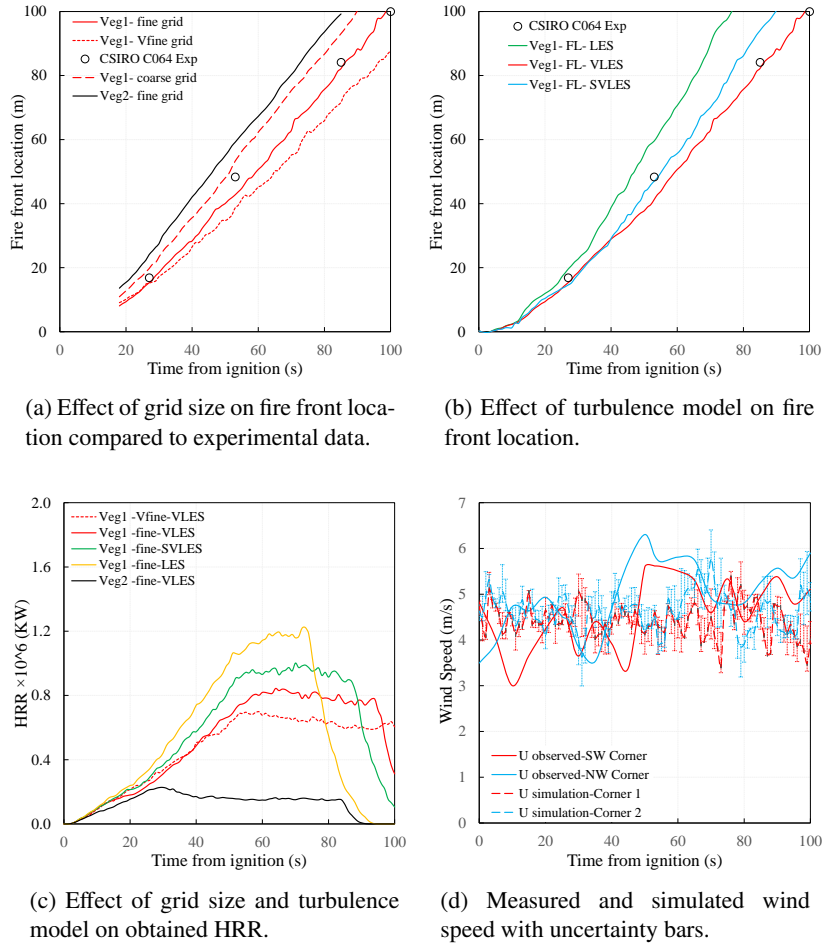
Figure 1: Model Domain.

A grid sensitivity analysis was performed to assess the sensitivity of the estimated *RoS* of the fire to the grid size. Three different grid resolutions of  $1 \times 1 \times 1 \text{ m}^3$  (coarse),  $0.5 \times 0.5 \times 0.5 \text{ m}^3$  (fine), and  $0.25 \times 0.25 \times 0.25 \text{ m}^3$  (very fine) were tested. Results are presented in Fig. 2a, which show a sensitivity to grid size ranging from 6% to 11% across different configurations of wind, vegetation height, and terrain slope. However, despite the benefits in accuracy, the use of the very fine grid is restricted by practical considerations such as the computational requirements (memory, processing) due to the large vegetation area.

Thus, a  $0.5 \times 0.5 \times 0.5 \text{ m}^3$  grid is used for the vegetation domain and its immediate surrounding ( $\pm 4 \text{ m}$ ) up to a height of 44 m. Upstream and downstream areas of the fine grid up to a distance of  $\pm 30 \text{ m}$  are discretised using a  $1 \times 1 \times 1 \text{ m}^3$  grid, while the rest of the domain is discretised using a  $2 \times 2 \times 1 \text{ m}^3$  grid up to a height of 60 m. A total of 16,687,680 grid cells are used, parallelised on 32 CPU cores. Fig. 2a compares the fire front location in the C064 experiment with estimations provided by three grid sizes,  $1 \times 1 \times 1 \text{ m}^3$  (coarse),  $0.5 \times 0.5 \times 0.5 \text{ m}^3$  (fine), and  $0.25 \times 0.25 \times 0.25 \text{ m}^3$  (very fine), and two vegetation types.

## 2.7. Model reliability and simulation scenarios

To examine the effect of the turbulence model on simulation accuracy and computational efficiency, three variants are compared: (i) LES, (ii) VLES, and (iii) simple VLES (SVLES). They offer different levels of physical detail and accuracy. Fig. 2 provides the fire front location and combustion heat release rate (HRR) for all grid sizes and turbulence models. As depicted in Fig. 2a, the estimated fire locations for the first vegetation model show agreement with the experimental measurements. The model shows grid dependence for both vegetation types, with the *RoS* decreasing as grid size decreases to 0.25 m. Further reduction in grid size did not affect the *RoS*. However, due to the high



**Figure 2:** Model sensitivity and validation analysis.

computational cost of using 0.25 m grid size, the 0.5 m grid size is adopted for the main simulations in this study to balance accuracy and computational cost for the large number of simulations carried out. Hence, the fuel parameters are adjusted to reproduce the CSIRO experiments while using a 0.5 m grid size. Fig. 2b shows the front location in the C064 experiment for the first vegetation model (Fir) and different turbulence models. Both LES and SVLES predict a faster RoS than the experiments, while VLES provides results that align with experimental data. The HRR for different grid sizes and turbulence models is depicted in Fig. 2c. Like RoS, it is sensitive to the grid size and turbulence model, though the variation is mostly limited to a maximum of 10 % for grid sizes of 0.25 m and 0.5 m. A comparison between the measured and simulated wind fields at field corners and 2 m AGL is provided in Fig. 2d. Comparing the measured and simulated wind signals indicates that, despite the stochasticity involved in simulating the wind signal, the error bounds suggest that the model predicts the average behaviour of the actual wind field within an acceptable range and captures the variability of velocity over time. Additionally, it should be mentioned that the simulation study presented here is generally idealized. It is important to keep in mind that atmospheric conditions such as wind velocity vary rapidly over time in field experiments, affecting fire propagation dynamics. These cannot be reproduced accurately in simulations, which are carried out under more uniform and controlled conditions. Differences between measurements and model predictions are partially due to this.

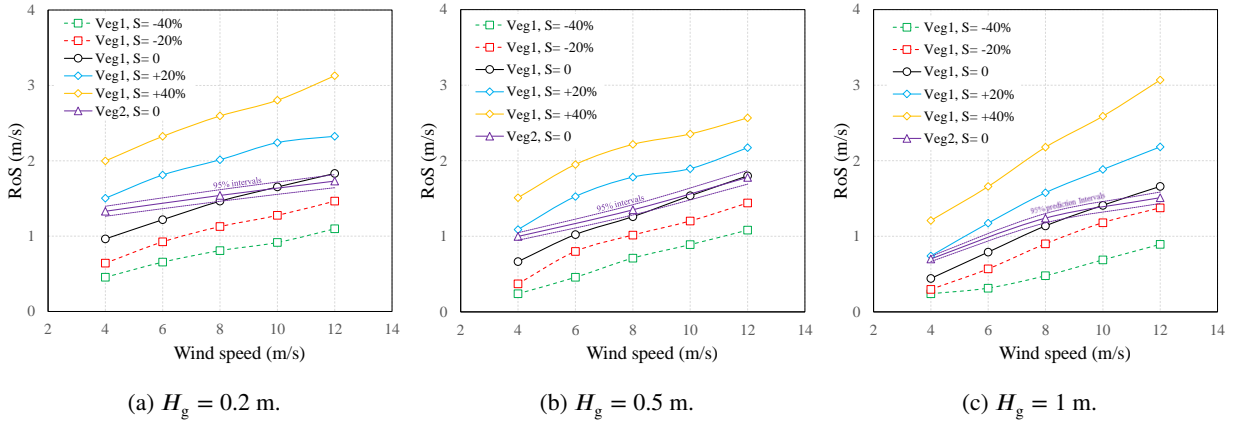
Comparing the simulated and measured fire contours on the  $x$ - $y$  plane for the CSIRO C064 experiment (not shown here), it is worth mentioning that, even though the simulated front location aligns with experimental data, the width and flanks of the fire front are less accurately estimated. This is mainly due to the smaller size of the flank flames compared to the front flames, requiring a finer grid to capture their dynamics accurately. To mitigate this and minimise



**Table 3**

Overview of the simulated geometric and physical properties of the vegetation, terrain, and wind

Vegetation model	Vegetation height (m)	Terrain slope ( $S$ )	Wind speed (m/s)	no. simulations
$Veg_1$	0.2, 0.5, 1	-40 % (-21.8°), -20 % (-11.31°), 0 % (0°), 20 % (11.31°), 40 % (21.8°)	4, 6, 8, 10, 12	75
$Veg_2$	0.2, 0.5, 1	0 % (0°)	4, 8, 12	9

**Figure 3:** Quasi-steady RoS at different wind speeds and slopes.

the effect of the length of the ignition line on the estimated RoS, it is kept the same as the width of the field (200 m) for all simulations.

A total of 84 simulations are performed, with 75 of them using the  $Veg_1$  and the remaining 9 using the  $Veg_2$  vegetation models. Three vegetation heights (0.2 m, 0.5 m, 1 m), various terrain slopes  $S \in [-21.8^\circ, 21.8^\circ]$ , and a range of wind speeds  $u_{10} \in [4, 12]$  m/s are considered in the simulations. A summary of the simulated scenarios is provided in Table 3. Note that  $Veg_2$  is only simulated on flat terrain to limit the total number of simulations.

### 3. Results and analysis

#### 3.1. RoS as a function of wind speed

The RoS is inherently dynamic and oscillatory, which is problematic for comparing different situations. Hence, a quasi-steady RoS is calculated for each scenario, which represents the average value. The quasi-steady RoS is considered to be the slope of the linear regression function fitted to the fire front locations. However, determining the fire front location can be challenging. Therefore, the fire front location at each time instance is identified as the front point along the centerline of the field where the temperature exceeds 400°C measured 25 cm AGL. This height is chosen to minimise the pulsating effect of the flames on the measurements.

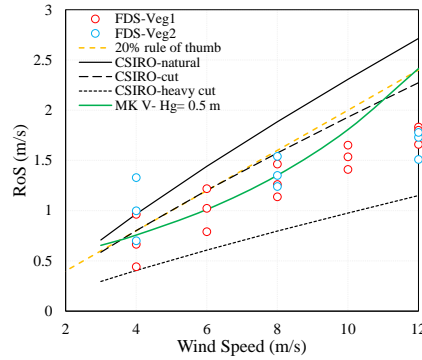
Fig. 3 shows the calculated quasi-steady RoS as a function of the wind speed ( $u_{10}$ ) for different terrain slopes ( $S$ ) and for three vegetation heights ( $H_g$ ). Generally speaking, it can be observed that higher wind speed leads to higher RoS regardless of the vegetation height and the bulk density. This increase mostly follows a linear trend, where the fuel with lower bulk density ( $Veg_2$ ) results in higher RoS than the fuel with higher bulk density ( $Veg_1$ ) on horizontal terrain, except for very high wind speeds  $u_{10} > 10$  m/s. Although each simulation was performed only once due to the associated high computational cost—making statistical uncertainty analysis infeasible, the 95 % intervals are added to the  $Veg_2$  values to highlight the difference between the estimated RoS of  $Veg_1$  and  $Veg_2$ .

Considering the combined effect of bulk density and wind speed, the RoS is observed to be higher at moderate wind speeds (4–8 m/s) for fuels with lower bulk density ( $Veg_2$ ), regardless of vegetation height. However, as the wind speed increases, the RoS of fuel with higher bulk density experiences a sharper increase and ultimately exceeds that of fuels with lower densities at high wind velocities ( $u_{10} > 10$  m/s).

For a constant bulk density, it can be observed that the RoS for shorter grass ( $H_g = 0.2$  m) tends to be higher than that for taller grasses ( $H_g = 0.5$  m and  $H_g = 1$  m), with the RoS decreasing for increasing grass heights. This

is consistent with the results in [27]. However, under high wind speeds  $u_{10} > 10$  m/s and steep slope conditions, the combined effects of vegetation height and terrain can be reversed, resulting in a higher RoS for taller grass ( $H_g = 1$  m), even with constant bulk density (see Figs. 3b and 3c).

Fig. 4 shows the quasi-steady RoS against wind speed over horizontal terrain obtained using FDS and three empirical models: the CSIRO [7], the McArthur mark V [34], and the 20% rule of thumb [10] models. All predictions are under fully cured conditions.



**Figure 4:** Quasy-steady RoS against wind speed over horizontal terrain obtained using FDS, CSIRO model, McArthur V model, and 20 % rule of thumb.

Since McArthur's model includes the effect of the fuel load, the equivalent fuel load for  $Veg_1$  with  $H_g = 50$  cm is utilised to calculate the corresponding RoS.

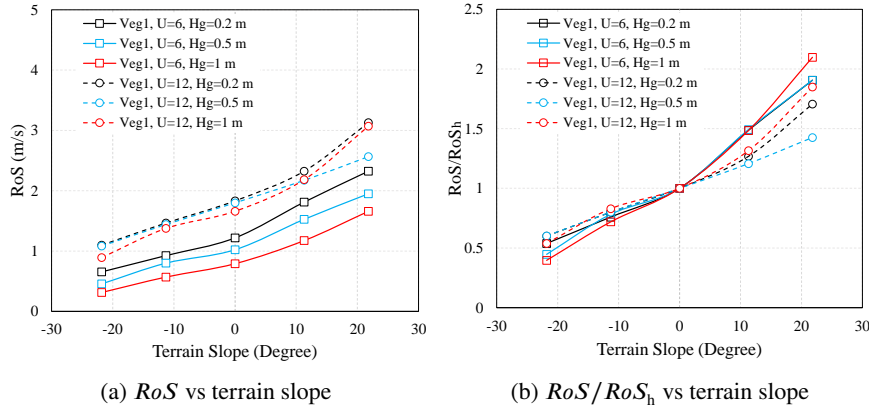
As expected, Fig. 4 shows that all models predict the RoS to increase with the wind speed, although the predicted values differ. These differences become larger for increasing wind speeds, with the FDS simulations clearly underestimating the empirical models' predictions. Notably, the CSIRO and the 20 % rule of thumb models show linear trends, while the influence of increasing wind speed on the RoS increases for the McArthur mark V model and decreases for FDS.

### 3.2. RoS as a function of terrain slope

As described in Section 2.7, the effect of bulk density on fire behaviour is studied solely on horizontal terrain. Therefore, this section focuses on the first fuel type ( $Veg_1$ ). Fig. 5a shows the quasi-steady RoS against terrain slope obtained by FDS for three vegetation heights and two wind speeds. Unsurprisingly, the RoS is higher for the higher wind speed, and increases monotonically with increasing slopes. Negative and positive slopes mean downhill and uphill propagation, respectively. As before, the RoS is higher for shorter grasses, which is more evident for  $u_{10} = 6$  m/s. This trend seemingly diminishes as wind speed increases. For  $u_{10} = 12$  m/s and small slopes, the influence of the grass height on the RoS is minimal.

Fig. 5b presents the ratio of the RoS at each terrain slope to the corresponding RoS over horizontal terrain ( $RoS/RoS_h$ ) for two wind speeds and three vegetation heights. An increase in the slope of the lines when  $S > |11.8^\circ|$  indicates that the influence of the terrain slope on the  $RoS/RoS_h$  ratio is more pronounced for larger magnitudes of the slope (whether uphill or downhill). Besides, it could be observed that the effect of the terrain slope on the  $RoS/RoS_h$  ratio is more pronounced at lower wind speeds: it increases faster uphill and decreases faster downhill. Under such conditions, fire propagation is primarily influenced by buoyancy forces, which leads to a plume-dominated propagation. This effect tends to diminish at higher wind speeds, as the fire propagation is then driven by the wind. Additionally, it can be concluded that the influence of terrain slope on accelerating the RoS on steeper terrain relative to the RoS on horizontal terrain ( $RoS/RoS_h$ ) is more significant in cases with tall vegetation ( $H_g = 1$  m), signifying the combined effect of vegetation and terrain slope on free propagation dynamics.

Fig. 6 shows the quasi-steady RoS against terrain slope obtained from FDS simulations (red and orange markers) for vegetation height  $H_g = 0.2$  m, and wind speeds  $u_{10} = 6$  m/s and  $u_{10} = 12$  m/s. These estimations are compared against those obtained from simulations in [20] (blue markers), and from the CSIRO [7] and the McArthur mark V [34] empirical models. The latter is corrected for slope effect using the correction factors in [34] as adopted in [20] and

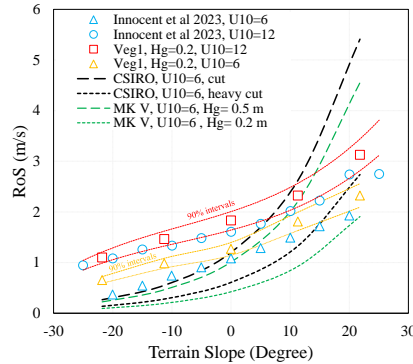


**Figure 5:** RoS and  $RoS/RoS_h$  ratio against terrain slopes for two wind speeds and three vegetation heights.

shown below:

$$RoS = RoS_h e^{0.069S} \quad (10)$$

where  $RoS$  is the forward RoS in km/h, and  $S$  is the slope of the terrain in degrees.



**Figure 6:** RoS against terrain slope from simulations (with 90% intervals) and from empirical models.

It may be inferred that,  $RoS$  from simulations in [20] with  $u_{10} = 12$  m/s mostly fall within the 90 % intervals of our simulations, simulations carried out with  $u_{10} = 6$  m/s show more significant differences.

Generally, the differences observed between our simulations and those in [20] are attributed to the use of different fuel models, wind models, grid sizes, and the longer 200-m ignition lines used in our study versus the 60-m ignition lines in [20]. Additionally, Innocent et al. [20] employed the FDS-BFM and a two-step thermal degradation model neglecting the exothermic char oxidation, while our study uses the FDS-LPM and a three-step reaction model. Finally, regarding the notable difference in results for  $u_2 = 6$  m/s, Innocente et al. [20] calibrated their model to reproduce Cheney [6] results with  $u_{10} = 6$  m/s. In contrast, our model was calibrated to match Cheney's results at the average wind velocity measured in the field, equal to  $u_2 = 4.6$  m/s (which corresponds to  $u_{10} \approx 5.75$  m/s), following the approach of Vanella et al. [46]. Given these differences in modelling and calibration approaches, the higher discrepancy at lower wind velocities is not unexpected.

Considering the agreement between simulations and empirical models, it is evident that both CSIRO-heavy cut and MK-V with  $H_g=0.2$  models show underestimated RoS values compared to our simulated values. However, the estimated values from the CSIRO-cut model are closer to those from the simulations within the range of  $S \in (-10, 10)$ .

### 3.3. RoS as a function of vegetation characteristics

In previous sections, we examined the combined effects of wind speed, terrain slope, and vegetation height on the RoS of the fire. Simulation results indicated that, generally, higher vegetation leads to lower RoS for constant

**Table 4**

Average and range of measured environmental and fire variables in [11] and [13].

Ref.	Vegetation		$T$ (°C)	$R_H$ (%)	$u_{10}$ (m/s)	$W$ (kg.m <sup>-2</sup> )	$H_g$ (m)	$\rho_b$ (kg.m <sup>-3</sup> )	$M$ (%)	$RoS$ (m/s)	$RoS/u_{10}$	$I$ (KW.m <sup>-1</sup> )
[11]	Wheat	mean	30.2	21.6	8.0	0.42	0.37	1.6	7.5	1.5	0.19	12131
		[range]	[24.8-38]	[13.6-31.7]	[5.2-10.8]	[0.32-0.53]	[0.08-0.83]	[0.73-4.04]	[5.4-11.6]	[0.7-2.8]	[0.08-0.38]	[3858-27987]
[13]	Grass	mean	25.7	27.8	5.3	0.49	0.39	1.38	7.96	1.0	0.18	7868
		[range]	[16-33]	[6-59]	[2-13]	[0.17-1.05]	[0.16-0.93]	[0.43-3.0]	[3.5-12.6]	[0.2-2.5]	[0.1-0.3]	[1260-18703]

bulk density ( $\rho_b$ ). This is in agreement with [27]. However, in some cases under high wind speeds and steep upslope conditions, the RoS is higher in taller grasses.

The practice of mowing grasses is widely adopted as a management strategy to control and slow down the wildfire spread in grasslands around the world, particularly favoured by the Australian fire authorities. The rationale is that cutting the grass would result in a less intense and slower propagating fire [12]. This is supported by experimental burns in [5] and [11]. In contrast, some studies such as [13] and [27] suggest that the RoS decreases with increasing vegetation height, and therefore the current practice of mowing grasses should be reconsidered. In this section, we address this question by analysing the RoS estimated from experiments and simulations under varying wind speeds over horizontal terrain. The aim is to explore how vegetation height and bulk density affect the RoS of the fire.

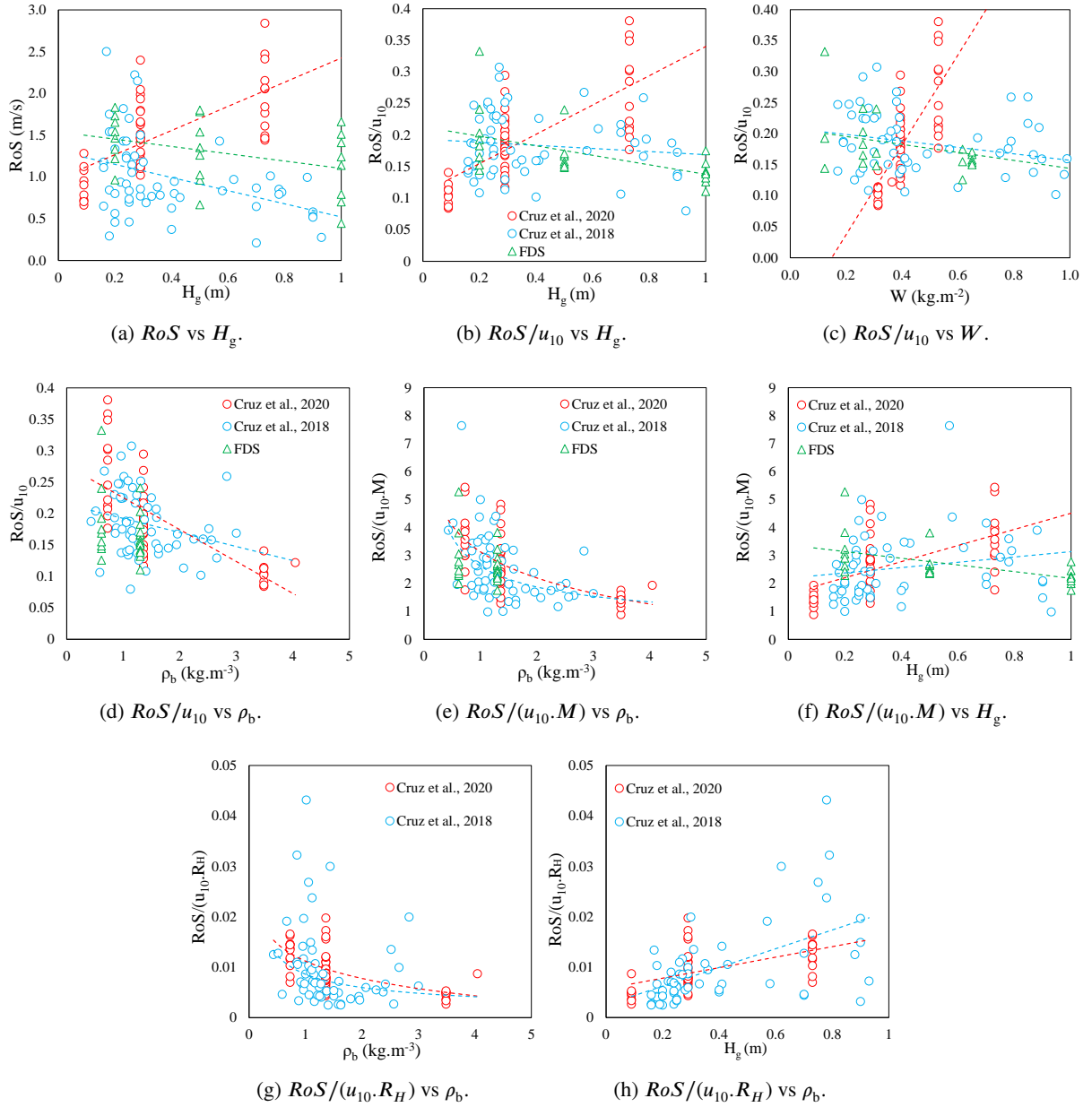
Two studies are selected with seemingly contradictory results regarding the relationship between fuel height ( $H_g$ ) and RoS of fire: grassland fire experiments in [13], and wheatland fire experiments in [11]. The aim is to investigate the sources of the contradictory results reported. Both studies' environmental and fuel conditions are summarised in Table 4. They are both performed in Australian territories under comparable temperature ( $T$ ), relative humidity ( $R_H$ ), vegetation height ( $H_g$ ), and moisture content ( $M$ ) conditions. It is important to note that the bulk density reported in [11] refers specifically to the standing fuel, while the burning experiments are carried out in the presence of matted fuel on the ground. However, for fire intensity calculations, the total fuel load includes 80 % of matted fuel load as well. Consequently, we calculate bulk density as the consumed fuel divided by the fuel height. While both the standing fuel bulk density and this calculated bulk density yield similar trends and insights, we opted to use the latter for better visualisation and to maintain consistency with other calculations such as fire intensity and with [13].

The average RoS and intensity ( $I$ ) from wheatland experiments in [11] ( $RoS_{avg} = 1.5$  m/s) are higher than those from grassland experiments in [13] ( $RoS_{avg} = 1$  m/s), which are seemingly due to the higher average wind speed during the wheatland experiments. However, experiments diverge primarily in their findings concerning the correlation between  $RoS$  and  $H_g$ . In the wheatland experiments, they are positively correlated with a sample Pearson correlation coefficient  $r = 0.672$ , in agreement with the conclusions made by Cheney et al. [5]. In contrast, the grassland experiments show a moderately negative correlation, with  $r = -0.39$  and  $p_r = 0.0024$ , where a  $p_r$ -value smaller than 0.05 denotes statistical significance. To reconcile these conflicting outcomes, Cruz et al. [12] state that the positive correlation between  $RoS$  and  $H_g$  is sufficiently confirmed experimentally. Consequently, they attributed the grassland study's contradictory conclusion to structural differences in the grasslands, which likely obscured the effect of  $H_g$  and led to misleading conclusions. Instead, Sutherland et al. [43] suggested that the differences are mainly due to different propagation modes: while the wheatland experiments primarily represented wind-driven fires, the grassland experiments were predominantly buoyancy-driven. They concluded that higher  $H_g$  leads to higher  $RoS$ .

Fig. 7 shows experimental data from [13] and [11] alongside data obtained from our FDS simulations. Properties studied includes fire  $RoS$ ,  $H_g$ , bulk density ( $\rho_b$ ), and fuel load ( $W$ ). We acknowledge the limitations—e.g., the infeasibility of calculating the statistical significance of correlation—and generalisation problems of using the sample Pearson correlation coefficient ( $r$ ) when comparing an explanatory variable (such as  $\rho_b$  or  $H_g$ ) with finite distinct values but multiple responses (i.e. RoS). This issue applies to our combined dataset, which includes both our simulations and data from [11]. However, given that part of the dataset is inherently grouped ([11]), while the other half remains ungrouped ([13]), alternative approaches are limited. Therefore, we analyse the data primarily utilising correlations and observable trends. Nonetheless, additional statistical methods such as ANOVA are employed to determine significance levels and provide a clearer overview of trends in the data from [11].

Fig. 7a shows the relationship between  $RoS$  and  $H_g$ , including experimental data from [13] and [11] alongside results from FDS simulations carried out in this study. The positive correlation in [11] ( $r = 0.67$ , with ANOVA yielding  $F = 22.03$  and  $p_F < 0.0001$ ) contrasts with the negative correlation from [13] ( $r = -0.4$ ,  $p_r = 0.0024$ ) and from our FDS simulations. A linear fit is applied for visualisation purposes only, acknowledging that it might not be the

### To Cut or Not to Cut



**Figure 7:** Distribution and trend of data including fire  $RoS$ , fuel height ( $H_g$ ), fuel load ( $W$ ), fuel moisture ( $M$ ), bulk density ( $\rho_b$ ), relative humidity ( $R_H$ ) and wind speed 10 m AGL ( $u_{10}$ ) for grassland experiments [13] (2018, blue markers), wheatland experiments [11] (2020, red markers), and our FDS simulations (2024, green markers).

best fit for all datasets. Notably, the use of the  $RoS$  may be misleading, as it is influenced by the ambient wind speed. Hence, using the  $RoS/u_{10}$  ratio provides a more reliable representation of the data, as depicted in Fig. 7b. A strong positive correlation between  $RoS/u_{10}$  and  $H_g$  for the wheatlands experiments is evident ( $r = 0.737$ , with  $F = 29.14$  and  $p_F < 0.0001$ ), while the significant negative correlation in grassland experiments no longer exists. This suggests that there is no  $(RoS/u_{10})-H_g$  correlation for grassland experiments in [13] or in our FDS simulations. A similar behaviour is observed for  $(RoS/u_{10})-W$  in Fig. 7c, suggesting that neither  $H_g$  nor  $W$  can sufficiently characterise the fire behaviour relationships and similarities in the two types of fuels.

However, when plotting the non-dimensional  $RoS/u_{10}$  ratio against  $\rho_b$ , a moderately weak yet statistically significant negative correlation emerges ( $r = -0.27$ ,  $p_r = 0.044$ ) for grassland data, whereas a stronger negative

correlation ( $r = -0.7$ , with  $F = 22.14$  and  $p_F < 0.0001$ ) is observed in the wheatland experiments (refer to Fig. 7d). It is generally consistent with the findings in [4] that increasing  $\rho_b$  leads to decreasing  $RoS$  for no-wind conditions. Here, this is still valid in the presence of wind.

In an effort to account for the effect of fuel moisture ( $M$ ), Fig. 7e shows  $RoS/(u_{10}M)$  against  $\rho_b$ , where a moderately strong negative correlation is observed ( $r = -0.57$ , with  $F = 15.06$ ,  $p_F < 0.001$ ) for wheatland fire and a moderately strong negative correlation ( $r = -0.42$ ,  $p_r < 0.001$ ) is observed for grassland fire dataset.

An intriguing behaviour is found when plotting  $RoS/(u_{10}M)$  against  $H_g$  in Fig. 7f. For both experimental datasets, a moderately strong positive trend is found ( $r = 0.59$ , with  $F = 29.14$ ,  $p_F < 0.0001$ ) for wheatland experiments, and a weakly positive though statistically insignificant correlation is found ( $r = 0.19$ ,  $p_r > 0.05$ ) for grassland experiments [13]. Despite the statistical insignificance of the latter trend, this highlights a practically consistent relationship between  $RoS$  and  $H_g$  when corrected for wind speed and moisture, reconciling the seemingly contradicting findings in [11] and [13]. This highlights the influence of additional parameters on the seemingly contradictory results of two experiments.

A closer examination of the environmental conditions in wheatland and grassland experiments suggests that atmospheric relative humidity ( $R_H$ ) could also be a key factor contributing to these discrepancies. Generally, grassland experiments are carried out under a wider range of atmospheric humidity values with a relatively higher average. The relative  $RoS$  values corrected for relative humidity are plotted as  $RoS/(u_{10}R_H)$  against  $\rho_b$  and  $H_g$  in Figs. 7g and 7h, respectively. Despite the seemingly similar trend of  $RoS/(u_{10}R_H)$  values with those of  $RoS/(u_{10}M)$ , the strength and statistical significance of trends are noticeably different. Fig. 7g indicates a moderately strong negative correlation for wheatland ( $r = -0.62$ ,  $F = 14.13$ ,  $p_F < 0.001$ ) and a moderately weak statistically insignificant negative correlation for grassland fire experiments ( $r = -0.22$ ,  $p_r > 0.05$ ). Finally, a moderately strong and statistically significant correlation is found between  $RoS/(u_{10}R_H)$  and  $H_g$ , with  $r = 0.55$  and  $p_r < 0.001$  for grassland fires, and  $r = 0.57$ ,  $F = 14.13$ , and  $p_F < 0.001$  for wheatland experiments. This highlights the major effects of environmental parameters on the contradiction observed in grassland experiments in [13]. It is acknowledged that the two bulk densities considered for FDS simulations may be insufficient to confidently establish a trend line. As a result, no trend is discussed for the FDS simulations presented in Figs. 7d, 7e, and 7g. Nevertheless, we conduct an in-depth analysis of the FDS results to evaluate its capability in correctly replicating the natural vegetation's height effect on its response to wildfires.

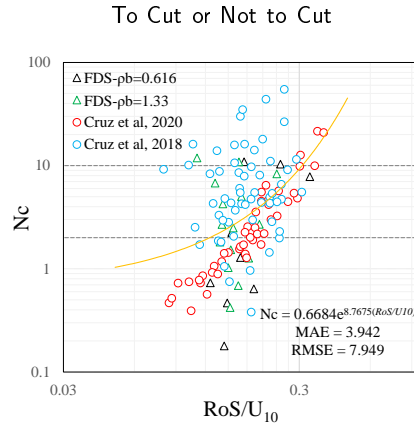
To summarise, a positive trend is found for  $RoS/u_{10}$  ratio with vegetation height ( $H_g$ ) when corrected for moisture ( $M$ ) or atmospheric humidity ( $R_H$ ). Nevertheless, the slopes of the trends in both datasets exhibit distinct magnitudes, reflecting the effect of various environmental parameters that affect fire behaviour, such as atmospheric humidity and temperature, bulk density, and the surface-area-to-volume ratio of the fuel which can even lead to contradictory and unexpected results which highlight the importance of further investigations.

However, contrary to experimental results, the output  $RoS/u_{10}$  of FDS simulations indicate a weak negative correlation even when corrected for moisture (Fig. 7f) which initially appears counter-intuitive. Section 4.1 delves into the potential reasons behind this discrepancy, and whether the FDS can accurately reproduce the correctly reproduce fire dynamics at wildland fires.

### 3.4. $RoS$ as a function of propagation mode

Despite the diverse range of burning conditions, fuel types, vegetation heights, terrain conditions, and other atmospheric parameters like moisture, a unifying perspective can be achieved by examining their collective impact on fire intensity ( $I$ ), and consequently, on the equilibrium of buoyancy and inertial forces within the fire front [30]. Therefore, Fig. 8 presents the  $RoS/u_{10}$  ratio against  $N_c$  (Byram number) for all data of horizontal terrain experiments and simulations used in this study, along with the best fit across the datasets. The best fit was achieved when combining experimental data with simulations performed on horizontal terrain, resulting in  $MAE = 3.942$  and  $RMSE = 7.949$ . This represents an improvement of 5–10% compared to using solely simulation data or solely experimental data. For wind-driven fires, particularly for  $N_c < 1$ , a saturation of data points is observed in the range  $RoS/u_{10} \in (0.08, 0.20)$  for all datasets, converging towards a state independent of  $N_c$ . This agrees with Morvan et al. [30], who suggested that fire propagation in this state is solely a function of  $M$ . However, a broad range of  $RoS/u_{10}$  values can be observed for larger values of  $N_c$ , increasing as  $N_c$  increases.





**Figure 8:**  $RoS/u_{10}$  ratio against the Byram number ( $N_C$ ) for the combined dataset of experimental results and horizontal terrain simulations. Dashed lines show  $N_C = 2$  and  $N_C = 10$ , where the fire propagation regime is wind-driven for  $N_C < 2$  and plume-dominated for  $N_C > 10$ . The solid orange line represents the best linear fit for the whole dataset.

## 4. Discussion

The previous section provided insight into the simulation results and the reanalysis of experimental datasets. However, questions were raised or remained partially unanswered, as they require the consideration of additional factors that affect the fire, including the fire propagation mode. Consequently, this section aims to combine analysis conclusions and fire propagation modes to address two key questions:

- (i) Is FDS capable of accurately reproducing the effect of grass height on wildfire propagation as observed in experimental field studies?
- (ii) Does cutting grass effectively slow down the propagation of wildfires?

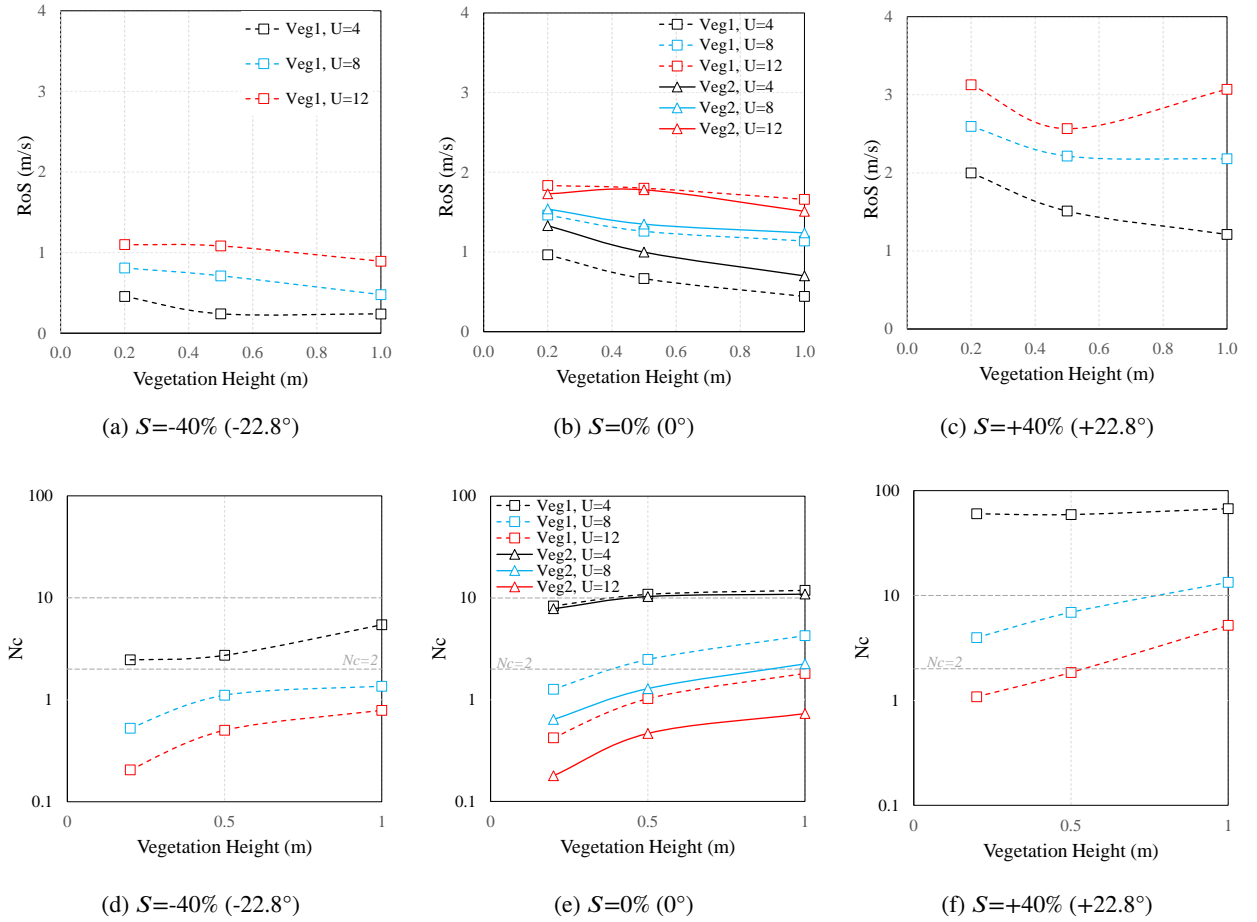
### 4.1. Are FDS results counter-intuitive?

The previous section revealed that FDS simulations describe a negative correlation between  $RoS/u_{10}$  and  $H_g$ . Similarly, Sutherland et al. [43, 27] reported a negative correlation between  $RoS$  and  $H_g$ , whereas experimental data found a positive correlation instead. However, it is crucial to highlight that all simulations in [43, 27] and the majority of our simulations assume constant bulk density ( $\rho_b$ ) across all  $H_g$ . In contrast, natural vegetation in the experimental datasets shows a decrease in  $\rho_b$  when increasing  $H_g$  with the Pearson correlation coefficient between  $H_g$  and  $\rho_b$  was found to be  $R_p = -0.62$  in [11] and  $R_p = -0.37$  in [13]. Cruz et al. [12] argues that changing  $H_g$  while keeping  $\rho_b$  constant—as in [43, 27] and in our simulations—solely changes the amount of fuel available for combustion, and that this does not suffice to explain the effect of these two fuel characteristics on the  $RoS$ . Besides, the representation of fuel as a homogeneous rigid layer that does not bend with the wind might have contributed to an incorrect estimation of drag coefficient and roughness length.

While we acknowledge the potential effect of vegetation representation and flexibility on the exerted drag and  $RoS$ , our analyses suggest that the disparity in  $\rho_b$  between simulated and natural vegetation significantly contributes to the observed differences in  $RoS/(u_{10}M)$  and  $RoS/(u_{10}R_H)$  with  $H_g$  correlations between FDS simulations and experimental observations. To further examine the effect of  $\rho_b$  on  $RoS$  and on the mode of propagation of the fire, Figs. 9a to 9c show the FDS predictions of how  $RoS$  and  $H_g$  are related at three wind speeds ( $u_{10}$ ), three terrain slopes ( $S$ ), and two  $\rho_b$ . The corresponding values of the Byram convective number ( $N_c$ ) for each of these cases are shown in Figs. 9d to 9f.

As previously mentioned, higher  $H_g$  with constant  $\rho_b$  leads to a reduction in the  $RoS$ , irrespective of the terrain slope or fire propagation mode. The only exception occurs for  $Veg_1$ ,  $S = 40\%$ , and  $u_{10} = 12$  m/s, in which case higher  $H_g$  between 0.5 m and 1 m results in higher  $RoS$ . This might be influenced by the shift in the fire propagation regime, as characterised by  $N_c$  (red line) crossing the  $N_c = 2$  threshold at  $H_g \approx 0.5$  m in Fig. 9f.

Thus, our findings are not in agreement with those in [43], which state that  $RoS$  increases with increasing  $H_g$  for constant  $\rho_b$  and wind-driven fires. However, it is important to note that their results are drawn from experiments with constant  $\rho_b$  and  $H_g < 0.2$  m, which is below the range considered in our simulations. Furthermore, the results obtained



**Figure 9:**  $RoS$  (a)–(c) and the corresponding Byram convective number ( $N_c$ ) (d)–(f) against vegetation height ( $H_g$ ) for two different vegetations,  $Veg_1$  and  $Veg_2$ , with respective bulk densities of  $\rho_b = 1.313 \text{ kg/m}^3$  and  $\rho_b = 0.616 \text{ kg/m}^3$ , at three wind speeds ( $u_{10}$ ) and three terrain slopes ( $S$ ). The  $N_c = 2$  and  $N_c = 10$  thresholds separate two fire propagation regimes, namely wind-driven ( $N_c < 2$ ) and plume-dominated ( $N_c > 10$ ).

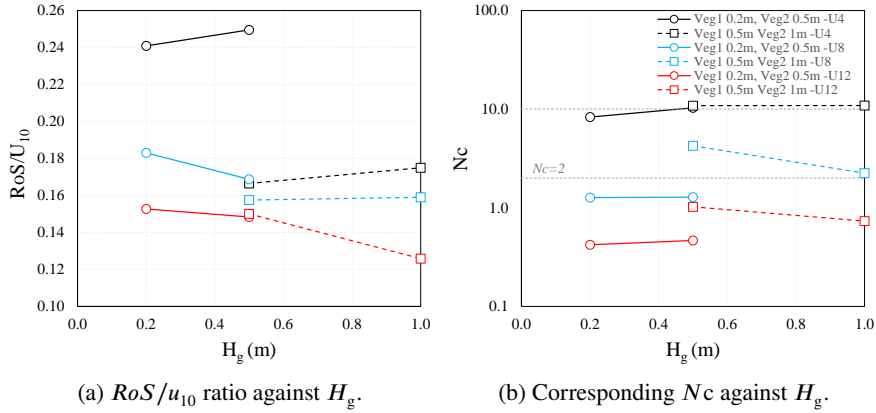
in our study using a wide range of wind speeds and terrain slopes do not confirm their observations, suggesting that factors beyond the fire propagation mode might have influenced the outcomes of their research.

Considering the effect of  $\rho_b$  on  $RoS$  and  $N_c$ , it can be observed in Figs. 9b and 9e that  $RoS$  for  $Veg_1$  (higher  $\rho_b$ ) is smaller than that of  $Veg_2$  (lower  $\rho_b$ ) for  $N_c \gg 2$  (plume-dominated and transitory regimes). This observation aligns with the conclusion in Section 3.3 that an increase in  $\rho_b$  leads to a reduction in  $RoS$ . Notably, this discrepancy tends to diminish as  $N_c \rightarrow 2$  (blue lines in Figs. 9b and 9e). However, for  $N_c$  well below the threshold of two (red lines, wind-driven fires), an interesting trend emerges:  $RoS$  for  $Veg_1$  ( $\rho_b = 1.33 \text{ kg/m}^3$ ) is slightly higher than that for  $Veg_2$  ( $\rho_b = 0.616 \text{ kg/m}^3$ ). This highlights a distinctive trend in wind-driven fires, which deviates from the general pattern observed in plume-dominated fires.

Fig. 10 depicts the  $RoS/u_{10}$  and corresponding  $N_c$  as a function of  $H_g$ , imitating the lower  $\rho_b$  characteristic seen in taller natural vegetation. For each case, the  $RoS/u_{10}$  of shorter grass is extracted from  $Veg_1$  with  $\rho_b = 1.33 \text{ kg/m}^3$ , while that of taller grass is extracted from  $Veg_2$  with  $\rho_b = 0.616 \text{ kg/m}^3$ . Consequently, for the solid black line,  $RoS$  at  $H_g = 0.2 \text{ m}$  comes from  $Veg_1$  and  $RoS$  at  $H_g = 0.5 \text{ m}$  corresponds to  $Veg_2$  with lower bulk density.

It is observed that for transitory and plume-dominated propagation modes ( $N_c > 2$ ), the  $RoS/u_{10}$  ratio increases with increasing  $H_g$  (simultaneously decreasing  $\rho_b$ ). However, for wind-driven fires ( $N_c < 2$ ), the  $RoS/u_{10}$  ratio decreases with increasing  $H_g$  (and simultaneously decreasing  $\rho_b$ ). This latter behaviour was not observed in the

wheatland experiments in [11], possibly due to the complexity and variability of experimental conditions and scattered data points, underscoring the need for further controlled experiments for wind-driven conditions.



**Figure 10:**  $RoS/u_{10}$  ratio and the corresponding Byram convective number ( $N_c$ ) against vegetation height ( $H_g$ ) for three wind speeds. For each line, the first point is extracted from  $Veg_1$  with  $\rho_b = 1.33 \text{ kg/m}^3$  and the second point from  $Veg_2$  with  $\rho_b = 0.616 \text{ kg/m}^3$ .

In summary, it has been found that FDS has the potential to realistically model the relationship between the  $RoS/u_{10}$  ratio and vegetation height ( $H_g$ ). This holds true provided that the model is correctly and accurately parametrised. Specifically, our analysis suggests that incorporating the natural decrease in bulk density ( $\rho_b$ ) as  $H_g$  increases—which is a characteristic observed in natural vegetation types—into the design of experiments, would allow FDS to better replicate the expected trend. Nevertheless, we do not dismiss the importance of conducting an uncertainty analysis and further sensitivity analysis to assess the statistical significance of the small increases or decreases observed. However, each simulation in this study was conducted only once due to the high computational cost of simulations and the negligible variations in quasi-steady RoS observed during our initial experimentations when repeating simulations.

#### 4.2. To cut or not to cut?

Mowing grasses is a widely accepted practice to curb the RoS of wildfires. While reducing  $H_g$  can indeed decrease fire intensity and therefore make the fire more manageable, it also makes it more prone to becoming wind-driven. Besides, our analyses suggest that the bulk density of the vegetation layer decreases ( $\rho_b \downarrow$ ) as the vegetation grows taller ( $H_g \uparrow$ ).

From the analyses in previous sections, it can be stated that increasing bulk density ( $\rho_b \uparrow$ ) reduces the rate of spread ( $RoS/u_{10} \downarrow$ ). Hence,  $RoS/u_{10}$  decreases with shorter grass ( $H_g \downarrow$ ). Therefore, mowing grasses may be an effective wildfire management strategy, as long as  $\rho_b$  increases as  $H_g$  decreases. However, it is important to note that this is valid for a constant moisture ( $M$ ) or  $R_H$ . However, variations in values of  $M$  or  $R_H$  may overshadow the effect of  $H_g$  as observed in [13] results.

The probability of wind-driven fire also increases when cutting the grass due to reduced fire intensity. Our simulations indicate that increasing bulk density ( $\rho_b \uparrow$ ) may lead to faster propagation ( $RoS/u_{10} \uparrow$ ) for wind-driven fires ( $N_c < 2$ ), though this behaviour was not observed in the wheatland experiments in [11]. Moreover, simulations in [43] suggest that taller vegetation may promote higher  $RoS$  for wind-driven fires, even if  $\rho_b$  is constant. However, further experimental research is required to determine if this effect can reduce the effectiveness or ultimately outweigh the benefits of reduced vegetation height.

## 5. Conclusions

The likelihood and intensity of wildfires have increased around the world due to global climate changes. Although grasslands cover less than 40 % of the Earth's surface, they encompass the majority of the burned area worldwide. This highlights the need for management strategies to confine the spread of wildfires in grasslands. One of these strategies is to cut or graze the grasses to limit the fire spread and intensity. However, this strategy has been challenged during the past few years based on results from a series of field experiments and numerical simulations.

This study combined contradicting field experiments with an extensive set of three-dimensional field-scale simulations to study the effect of vegetation height and vegetation bulk density on wildfire propagation dynamics under different atmospheric and terrain conditions. Results indicate that increasing terrain slope (uphill conditions) or wind speed leads to an increased rate of spread, irrespective of the vegetation height and bulk density.

As a key finding, this study identifies the critical role of bulk density in fire dynamics, showing that the rate of fire spread decreases with increasing vegetation bulk density. It was found that the bulk density of natural vegetation—at least those considered in this study—decreases as the vegetation grows taller. Thus, the relationship between vegetation height and fire RoS seems to be inherently linked to the variation in vegetation bulk density. In this study, we also aimed to establish general relationships between  $RoS/u_{10}$  and vegetation height ( $H_g$ ), acknowledging the challenges that arise from numerous environmental factors that can influence or even obscure the effects of vegetation height and bulk density. To account for these complexities, we introduced parameters such as  $RoS/(u_{10}M)$  and  $RoS/(u_{10}R_H)$  to adjust  $RoS/u_{10}$  for the specific environmental conditions of each experiment. These adjustments allowed us to uncover statistically significant correlations across both vegetation types, providing insight into fire spread dynamics.

Simulation results suggest that increasing the vegetation height, while maintaining bulk density constant, leads to a reduced RoS. Conversely, in natural vegetation, where the bulk density decreases with increasing height, the RoS increases for taller vegetation. This implies that reduced bulk density is a primary driver, as opposed to vegetation height alone. However, it must be noted that, even though the above-mentioned behaviours were observed in experimental data of both plume-dominated and wind-driven propagation modes, further simulations suggested potential reversals under wind-driven fire conditions. The lack of experimental data to validate the extreme behaviours observed in simulations emphasises the importance of designing and performing more field tests to support the investigation of fire dynamics. Such experimental studies are crucial for refining existing and projected grassland fire management strategies, especially for regions prone to high wind conditions.

Finally, we found that FDS has the potential to realistically reproduce the relationship between the  $RoS/u_{10}$  and vegetation height ( $H_g$ ), provided that it is properly parameterized. In particular, our analysis suggests that accounting for the natural decrease in bulk density ( $\rho_b$ ) as  $H_g$  increases—a pattern seen in real vegetation—would help FDS more accurately reproduce the expected trend. Yet, further experimentations and statistical analysis would be required to accurately analyse the model's sensitivity and responses.

## CRediT authorship contribution statement

**Mohammad Tavakol Sadrabadi:** Conceptualization, Methodology, Software, Writing, Investigation, Visualization. **Mauro Sebastián Innocente:** Conceptualization, Methodology, Software, Writing, Supervision, Resources, Visualization.

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